



# Investigation of shape effects of Cu-nanoparticle on heat transfer of MHD rotating flow over nonlinear stretching sheet



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Grashof number;  
Numerical consequence

**Abstract** In the current study, we focused on the shape effects of copper (Cu) nano-particles on heat transmission of three-dimensional magnetohydrodynamic (MHD) nano-fluid. A particular type of flow is considered, i.e., rotating flow over an exponentially stretching sheet. Significant actions of thermal radiation and Grashof number are also formulated to study. The modified Chebyshev wavelets method is introduced to examine the numerical solutions of the accomplished model. Error analysis and further comparison with existing results reflect the appropriateness of the proposed modification. Graphical behavior of dimensionless velocities, temperature, Nusselt number, and skin friction under the inspiration of several parameters is also presented. The attained solutions propose that the modification is beneficial and can be protracted to other highly nonlinear problems. The modified Chebyshev wavelets technique lowers computing time, according to the study. The research takes into account the form impacts of nano-particles. There are three kinds of nano-particles to examine (spherical, laminar, and cylinder). The behavior of various particles

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varies depending on the flow and temperature profile. For microparticles with a laminar structure, velocity and temperature have higher values.

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## 1. Introduction

Nowadays, the concept of nano-fluids has become essential and significant for almost every area of science and technology. This is due to the enormous applications of this field of science. It is directly applicable for many areas like in the study of engineering (electronics, industry, printing, devices linked to military) and medical-related sciences (surgery connected with the body, therapy of many dangerous diseases, medicine, food, and nutrition). It has found its applications in the study related to materials and their properties, in the field of computers (in particular, for the computers having unlimited powers), studying the nuclear-based reactor, etc. An overview of the literature related to fluid mechanics exhibits that many theories, laws, and models have been established so far. Basically, nano-fluids contain small particles (named nano-particles) in the base of fluids. Due to these specialized particles, fluid properties are altered and adjusted, especially conductivity (thermal) [1] and heat transmission processes. Choi [2] was the first to investigate the improvement of nano-sized particles in fluids. Later, Lee et al. [3] developed the idea of nano-particles and examined the fluid that surrounds the nanotube. Multiple investigators have begun researching in this area due to its numerous applicability in various engineering and sciences applications [4–11] (see Table 1).

Radiative fluids are essentially such a category (of fluids) that yield energy due to radiations settled in the system geometry. This imperative tradition for the circumstance of fluid is significant in countless industrial measures. This category of fluids is exclusively appropriate in those structures in which a particular type of fluid is baptized as “ambient” fluid. Moreover, this research is compelling for several categories of devices in specific areas of space science, which can function at a hefty temperature scale. Frequently, “Stefan approximations” of radiation are applied to different mathematical formulations. Micropolar perceptions, along with radiation enrichment, are offered by Ishak [6]. T. Hayat [12], a renowned mathematician, discussed the unsteady wonders of radiative and magnetic fluid in the same year. Shateyi et al. [13] proposed a method for calculating geometric significances in laminar and radiative flow. Khan [14] et al. put their findings on MHD nano-fluid characteristics and heat energy production perceptions into practice.

Learning MHD procedures and embedding this information on nano-fluid have become very effective and validate

for large regions of sciences such as optical modulators and wound behaviors. Moreover, theoretical study of this field provides the knowledge of special forces termed “Lorentz force.” This support exists because of the MHD study, which obliges the theme control the chilling (cooling) structures. To model this category of phenomena, the philosophy of Ohm’s laws and Maxwell are very significant. The modern thoughts concerning this range can be perceived using these refs. [12–20,22,23,24].

The flow of rotation types is a very significant phenomenon. They are highly functional in numerous areas of science like rotating type of machinery, the flow of magma in the mantle layer of earth, flow of anti-cyclonic, the process of centrifugal filtration, chemical and food procedures, and viscometers. A famous scientist Crane [21], fashioned his important investigation in the arena of flow (stretching). Fuhrer, Wang [22] contributed to study the rotational fluid fashioned by enlarging surface. Further, he cracked this demonstrated problem utilizing perturbation for a slight alteration in parameter  $k$ . Afterward, abundant reaches worked on these philosophies and verbalized countless models using rotational fluidic properties. Supplementary concepts can be perceived from specified refs. [23–26,27–48].

The literature review in the preceding paragraphs is only focused on the research of shape impacts of nano-particles on 3D rotating nano-fluid flow across a stretched surface with magnetic field and heat radiation. The numerical results of the proposed modeled physical issue are assessed using the modified Chebyshev wavelets technique, which is a numerical methodology. In the next part of the article, graphical analysis and parametric research are also given to investigate the physical implications of the modeled issue. The proposed method may be tweaked using polynomial and wavelet theory literature analysis.

## 2. Mathematical and geometrical analysis

Assume three-dimensional rotating and convective free flow of nano-fluid over an elastic sheet occupying  $z = 0$  plane. The sheet is continuously stretching with velocity whose components are  $U_w$  and  $V_w$  in  $x$  and  $y$ -axis directions, as shown in Fig. 1. Consider the flow is steady, laminar boundary layer, and incompressible. Further, consider the flow rotates with an angular velocity  $\Omega = \Omega_0 e^{-(x+y)/L}$ , here  $L$  and  $\Omega_0$  are reference length and average form of velocity, respectively.

**Table 1** Thermo-physical properties of blood and nano-particles [11].

Properties	$\beta(\text{K}^{-1})$	$\sigma(\Omega\text{m})$	$c_p(\text{J/kgK})$	$k(\text{W/mK})$	$\rho(\text{kg/m}^3)$
Water	$21 \times 10^{-5}$	0.05	4179	997	0.613
Copper (Cu)	$1.67 \times 10^{-5}$	$5.96 \times 10^7$	385	8933	401

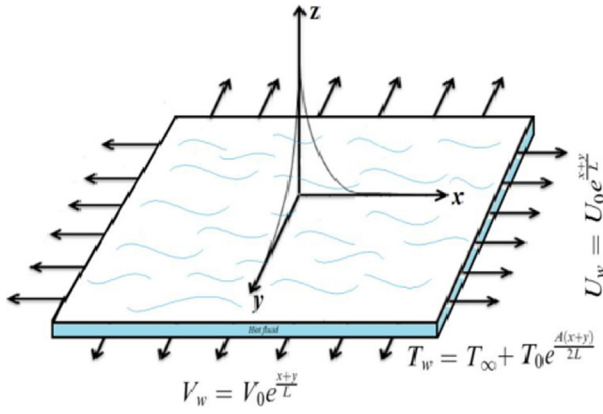


Fig. 1 Systematic diagram of the problem.

A magnetic field of constant nature is applied, and a chemical reaction of the first order is implemented to the system. It is also part of the assumption that all the other sources, such as Joule heating, which affects fluid temperature, are negligibly small. Utilizing the above assumptions, boundary layer approximations, we get the following mathematical model [34]:

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = 0, \quad (1)$$

$$u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} + w \frac{\partial u}{\partial z} - 2\Omega v = v_{nf} \frac{\partial^2 u}{\partial z^2} + \frac{(\rho\beta)_{nf}}{\rho_{nf}} g(T - T_\infty) - \frac{\sigma_{nf} B_0^2}{\rho_{nf}} u, \quad (2)$$

$$u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} + w \frac{\partial v}{\partial z} - 2\Omega u = v_{nf} \frac{\partial^2 v}{\partial z^2} + \frac{(\rho\beta)_{nf}}{\rho_{nf}} g(T - T_\infty) - \frac{\sigma_{nf} B_0^2}{\rho_{nf}} v, \quad (3)$$

$$\begin{aligned} & \frac{(\rho c_p)_{nf}}{(\rho c_p)_f} \left( u \frac{\partial T}{\partial x} + v \frac{\partial T}{\partial y} + w \frac{\partial T}{\partial z} \right) \\ & = \frac{k_{nf}}{(\rho c_p)_f} \frac{\partial^2 T}{\partial z^2} + \frac{\mu}{(\rho c_p)_f} \left( \left( \frac{\partial u}{\partial z} \right)^2 + \left( \frac{\partial v}{\partial z} \right)^2 \right) + \frac{1}{(\rho c_p)_f} \frac{\partial}{\partial z} q_r, \end{aligned} \quad (4)$$

The useful relations for nano-particles concepts are explained as [11]:

$$\begin{aligned} \frac{\mu_{nf}}{\mu_f} &= \frac{1}{(1-\phi)^{2.5}}, \quad \frac{\rho_{nf}}{\rho_f} = \left( 1 - \phi + \frac{\rho_s}{\rho_f} \phi \right), \\ \frac{\sigma_{nf}}{\sigma_f} &= \left( 1 + \frac{3(\sigma-1)\phi}{(\sigma+2) - (\sigma-1)\phi} \right), \\ \frac{(\rho c_p)_{nf}}{(\rho c_p)_f} &= \left( 1 - \phi + \frac{(\rho c_p)_s}{(\rho c_p)_f} \phi \right), \\ \frac{(\rho\beta)_{nf}}{(\rho\beta)_f} &= \left( 1 - \phi + \frac{(\rho\beta)_s}{(\rho\beta)_f} \phi \right), \\ \frac{k_{nf}}{k_f} &= \frac{k_s + (n-1)k_f - (n-1)(k_f - k_s)\phi}{k_s + (n-1)k_f + (k_f - k_s)\phi}, \quad \sigma = \frac{\sigma_s}{\sigma_f}. \end{aligned}$$

The suitable boundary conditions (BC's) are connected with the above Eqs. (1)–(4) are:

$$u = U_w, v = V_w, w = 0, T = T_w, \text{at } z = 0, \quad (5)$$

$$u = v = 0, T = T_\infty, \text{as } z \rightarrow \infty. \quad (6)$$

Further we have:

$$U_w = U e^{\frac{x+y}{2L}}, V_w = V e^{\frac{x+y}{2L}}, T_w = T_\infty + T_0 e^{\frac{x+y}{2L}}.$$

In above  $U, V$  and  $T_0$  are constants,  $T_\infty$  is ambient temperature. The last term of Eq. (4) signifies the influence of thermal radiation, which can be written as [11]:

$$\frac{\partial}{\partial z} q_r = -\frac{16}{3} \frac{\sigma^* T_\infty^3}{k^* (\rho c_p)_f} \frac{\partial^2 T}{\partial z^2}.$$

According to the problem, we have the following similarity transformations:

$$\begin{cases} u = U e^{\frac{x+y}{2L}} F'(\eta), v = U e^{\frac{x+y}{2L}} G'(\eta), \theta(\eta) = \frac{T - T_\infty}{T_0} e^{-\frac{A(x+y)}{2L}}, \\ w = -\sqrt{\frac{U}{2L}} e^{\frac{x+y}{2L}} \left[ F(\eta) + \eta F'(\eta) + G(\eta) + \eta G'(\eta) \right], \eta = z \sqrt{\frac{U}{2L}} e^{\frac{x+y}{2L}}. \end{cases} \quad (7)$$

We obtained the following set of equations by using the above transformation into the modelled problem (1)–(4):

$$\begin{aligned} \frac{v_{nf}}{v_f} F'''' + F''(F + G) - 2F'(F' + G') + 4\lambda G' + \frac{\rho_f}{\rho_{nf}} \frac{(\rho\beta)_{nf}}{(\rho\beta)_f} Gr\theta \\ = \frac{\rho_f}{\rho_{nf}} \frac{\sigma_{nf}}{\sigma_f} Ha^2 F', \end{aligned} \quad (8)$$

$$\begin{aligned} \frac{v_{nf}}{v_f} G'''' + G''(F + G) - 2G'(F' + G') - 4\lambda F' + \frac{\rho_f}{\rho_{nf}} \frac{(\rho\beta)_{nf}}{(\rho\beta)_f} Gr\theta \\ = \frac{\rho_f}{\rho_{nf}} \frac{\sigma_{nf}}{\sigma_f} Ha^2 G', \end{aligned} \quad (9)$$

$$\begin{aligned} \frac{1}{Pr} \left( \frac{k_{nf}}{k_f} + Rd \right) \theta'' + \frac{(\rho c_p)_{nf}}{(\rho c_p)_f} \left( (F + G)\theta' - A(F' + G')\theta \right) \\ + Ec_x Ec_y (F''^2 + G''^2) = 0. \end{aligned} \quad (10)$$

Boundary condition gets reduces to as:

$$\begin{cases} f(\eta) = g(\eta) = 0, f'(\eta) = 1, g'(\eta) = \alpha, \theta(\eta) = 1, \text{at } \eta = 0, \\ f'(\eta) = g'(\eta) = 0, \theta(\eta) = 0, \text{as } \eta \rightarrow \infty. \end{cases} \quad (11)$$

Important symbols and expressions are given in Table 2.

### 3. Wavelets and Chebyshev wavelets

The third kind Chebyshev wavelets (TKCW) with the influences  $k, p, j, M$  defined in [0,1] is given as below [35,36]:

$$\psi_{p,j}(\eta) = \begin{cases} 2^{\frac{k}{2}} V_j(2^k \eta - (2p-1)), & \frac{p-1}{2^{k-1}} \leq \eta \leq \frac{p}{2^{k-1}}, \\ 0 & \text{otherwise,} \end{cases}$$

where  $V_j(\eta) = \frac{1}{\sqrt{\pi}} V_j(\eta)$  is represents the  $j$ -th-order Chebyshev wavelets. Also  $j = 0, 1, \dots, M-1, p = 0, 1, \dots, 2^{k-1}$ . The recurrence relation of Chebyshev polynomials of the third kind are explained as:

$$V_0(\eta) = 1, V_1(\eta) = 2\eta - 1, V_j(\eta) = 2\eta V_{j-1}(\eta) - V_{j-2}(\eta), j \geq 2.$$

The weight function  $\Xi_p$  of Chebyshev wavelets are give below:

**Table 2** Explanation of symbols & expressions (nomenclature).

Name	Symbols&Expressions	Name	Symbols&Expressions
Kinematicviscosity	$\nu$	Thermal conductivity	$k$
Subscriptfornano – fluid	$nf$	Thermal expansion	$\beta$
Electricalconductivity	$\sigma$	Mass diffusion	$D$
Density	$\rho$	Concentration susceptibility	$c_s$
Magneticfieldcomponent	$B_0$	Thermal diffusion ratio	$k_T$
Temperature	$T$	Shapefactor	$n$
Heatcapacity	$c_p$	Volumefraction	$\phi$
Subscriptforbasefluid	$f$	Eckertnumbersinx	$Ec_x = \frac{v_f U^2}{\alpha c_p T_0} e^{\frac{2x}{L} - \frac{4x}{2L}}$
Rotation	$\lambda = \frac{\Omega L}{U} e^{-\frac{x+y}{L}}$	Eckertnumbersiny	$Ec_y = \frac{v_f U^2}{\alpha c_p T_0} e^{\frac{2y}{L} - \frac{4y}{2L}}$
Hartmannnumber	$Ha^2 = \frac{2\sigma_f L B_0^2}{U \rho} e^{-\frac{x+y}{L}}$	Stretchingratio	$\alpha = \frac{V}{U}$
Grashofnumber	$Gr = \frac{2L\beta_f g T_0}{U^2} e^{\frac{A(x+y)}{2L} - \frac{2(x+y)}{L}}$	Radiation	$Rd = \frac{16\sigma^* T_\infty^3}{3k k_f}$
Prandtlnumber	$Pr = \frac{k_f}{\nu_f (\rho c_p)_f}$	Skin Friction inx	$\sqrt{\frac{Re_x}{2}} C_{fx} = -e^{\frac{3(x+y)}{2L}} \frac{\rho_f \mu_{nf}}{\rho_{nf} \mu_f} F''(0)$
Nusselt number	$\frac{\sqrt{2}L}{x} \frac{Nu_x}{\sqrt{Re_x}} = -\left(\frac{k_{nf}}{k_f} + Rd\right) e^{\frac{x+y}{2L}} \theta'(0)$	Skin Friction iny	$\sqrt{\frac{Re_x}{2}} C_{fy} = -e^{\frac{3(x+y)}{2L}} \frac{\rho_f \mu_{nf}}{\rho_{nf} \mu_f} G''(0)$
Reynolds number	$Re = \frac{UL}{\nu}$	Angular velocity	$\Omega = \Omega_0 e^{-(x+y)/L}$

$$\Xi_p(\varrho) = \Xi(2^k \varrho - (2p_1 - 1)) = \sqrt{\frac{1 + (2^k \varrho - (2p_1 - 1))}{1 - (2^k \varrho - (2p_1 - 1))}}$$

With the concepts of TKCW, we can write as:

$$F(\varrho)|_{L^2(\mathbb{R},[0,1])} = \sum_{p=1}^{\infty} \sum_{j=0}^{\infty} \xi_{pj} \psi_{p,j}(\varrho) \Big|_{\xi_{pj} = \int_0^1 F(\varrho) \psi_{p,q}(\varrho) \Xi_p d\varrho = (F(\varrho), \psi_{p,q}(\varrho))_{L^2_{\Xi_p}[0,1]}}$$

The above-defined equation is truncated as:

$$F(\varrho) = \sum_{p=1}^{2^{k-1}} \sum_{q=0}^{M-1} \xi_{pq} \psi_{p,q}(\varrho) = \mathcal{M}^T \psi(\varrho),$$

where  $\psi(\varrho)$  and  $\mathcal{M}$  are specified in [35–38].

#### 4. Solution procedure

This section scrutinizes the methodology and application of the modified Chebyshev wavelets method to find the solution (8–10). This revised version has the following steps:

**Step 1.** First, consider (8–10):

$$\frac{\nu_{nf}}{\nu_f} F''' + F''(F + G) - 2F'(F' + G') + 4\lambda G' + \frac{\rho_f}{\rho_{nf}} \frac{(\rho\beta)_{nf}}{(\rho\beta)_f} Gr\theta = \frac{\rho_f}{\rho_{nf}} \frac{\sigma_{nf}}{\sigma_f} Ha^2 F', \tag{12}$$

$$\frac{\nu_{nf}}{\nu_f} G''' + G''(F + G) - 2G'(F' + G') - 4\lambda F' + \frac{\rho_f}{\rho_{nf}} \frac{(\rho\beta)_{nf}}{(\rho\beta)_f} Gr\theta = \frac{\rho_f}{\rho_{nf}} \frac{\sigma_{nf}}{\sigma_f} Ha^2 G', \tag{13}$$

$$\frac{1}{Pr} \left( \frac{\kappa_{nf}}{\kappa_f} + Rd \right) \theta'' + \frac{(\rho c_p)_{nf}}{(\rho c_p)_f} ((F + G)\theta' - A(F' + G')\theta) + Ec_x Ec_y (F'^2 + G'^2) = 0. \tag{14}$$

**Step 2.** The trial solution according to Chebyshev wavelets method for solving Eqs. (12)–(14) are given as:

$$\mathcal{F}(\aleph) = \mathbf{C}_l^T \psi(\varrho), \tag{15}$$

where  $\sum_{i=1}^{2^{k-1}} \sum_{j=0}^{M-1} \Delta_{ij}^q \mathcal{H}_{ij}(\aleph) = \mathbf{C}_l^T \psi(\varrho)$ ,  $q = l = 1, 2, 3$  for  $\mathcal{F}(\aleph) = (F(\aleph), G(\aleph), \theta(\aleph))$  respectively.

$$\mathbf{C}_l = \left[ \Delta_{1,k}^i \right]^T, k = 0, 1, 2, \dots$$

The above solution (15) can be rephrased as:

$$\mathcal{F}(\aleph) = \Lambda_l^T \chi(\varrho), \chi(\varrho) = [1, \varrho, \varrho^2, \varrho^3, \dots]^T. \tag{16}$$

where  $\Lambda_p, p = 1, 2, 3$  are explained in refs. [31,32]. Therefore; we further have

$$\mathcal{F}(\aleph) = \sum_{n=0}^M \zeta_n^1 \varrho^n. \tag{17}$$

**Step 3.** After inserting the trial solutions (22–24) into Eqs. (8)–(10), we got the following residuals for velocities and temperature:

$$\mathbf{R}_F = \frac{\nu_{nf}}{\nu_f} F''' + F''(F + G) - 2F'(F' + G') + 4\lambda G' + \frac{\rho_f}{\rho_{nf}} \frac{(\rho\beta)_{nf}}{(\rho\beta)_f} Gr\theta - \frac{\rho_f}{\rho_{nf}} \frac{\sigma_{nf}}{\sigma_f} Ha^2 F',$$

$$\mathbf{R}_G = \frac{\nu_{nf}}{\nu_f} G''' + G''(F + G) - 2G'(F' + G') - 4\lambda F' + \frac{\rho_f}{\rho_{nf}} \frac{(\rho\beta)_{nf}}{(\rho\beta)_f} Gr\theta - \frac{\rho_f}{\rho_{nf}} \frac{\sigma_{nf}}{\sigma_f} Ha^2 G',$$

$$\mathbf{R}_\theta = \frac{1}{Pr} \left( \frac{\kappa_{nf}}{\kappa_f} + Rd \right) \theta'' + \frac{(\rho c_p)_{nf}}{(\rho c_p)_f} ((F + G)\theta' - A(F' + G')\theta) + Ec_x Ec_y (F'^2 + G'^2).$$

**Step 4.** Now Galerkin method is applied in order to analyze the values of  $\zeta$ 's. We have:

$$\mathcal{H}_\ell^n = \int_0^{\ell_\infty} \mathbf{R}_\ell \frac{d}{dt_n} \ell(\mathbb{N}) d\mathcal{Q}, n = 1, 2, \dots, 2^{k-1}M - 3,$$

where  $\ell = (F, G, \theta)$ ,  $\mathcal{H}_\ell^n = (\mathbf{E}_F^n, \mathbf{E}_G^n, \mathbf{E}_\theta^n)$ ,  $\mathbf{R}_\ell = (\mathbf{R}_F, \mathbf{R}_G, \mathbf{R}_\theta)$  and  $l = 1, 2, 3$  for  $F, G$  and  $\theta$  respectively.

**Step 5.** The values of  $\Delta$ 's (unknowns) is achieved. Consequently, the values of  $\Delta$ 's accomplished by inserting the unknowns  $\Delta$ 's.

**5. Results and discussion**

To explore the physical features of both parametric and shape effects on velocity and temperature, numerical solution of the expressed model is offered in the earlier section, has been evaluated using a modified version of scheme modified Chebyshev wavelets method. Graphical analysis for shape and parametric effects on physical quantities of interest, temperature, and velocity are presented in Figs. 2–8. Furthermore, results and

discussion regarding Figs. 2–8 are also part of this portion of the paper. For simplicity purposes, we consider  $\varrho = \eta$ .

The “Nusselt number” and “skin friction” are explained in Figs. 2 and 3. It can be seen that skin friction coefficient both  $x$  and  $y$ -direction demonstrate the growing behavior against the Hartmann number. This energy is added to the system, and further fluid particles are more excited due to it. Therefore, it becomes a significant increase in temperature and velocity. Due to the variation in rotating parameter coefficient ( $x$ -directional skin friction) decreases gradually, and the reverse effect of rotating parameter achieved for  $y$ -directional “skin friction.” It has the dominant values for the case of the sphere shape of the nano-particle. Nusselt number ( $Nu_x$ ) upsurges due to the changes in  $Pr$ ,  $\alpha$ , and  $\beta$ .  $Nu_x$  decrease as enhancing the  $Ec_x$ . Similar results of the Nusselt number can be obtained for the Eckert number in  $y$ -direction.

Furthermore, Figs. 4–8 are graphical behavior analyses of velocity and temperature for distinct numerical values of parameters. Fig. 4(a) displays the effects of velocity

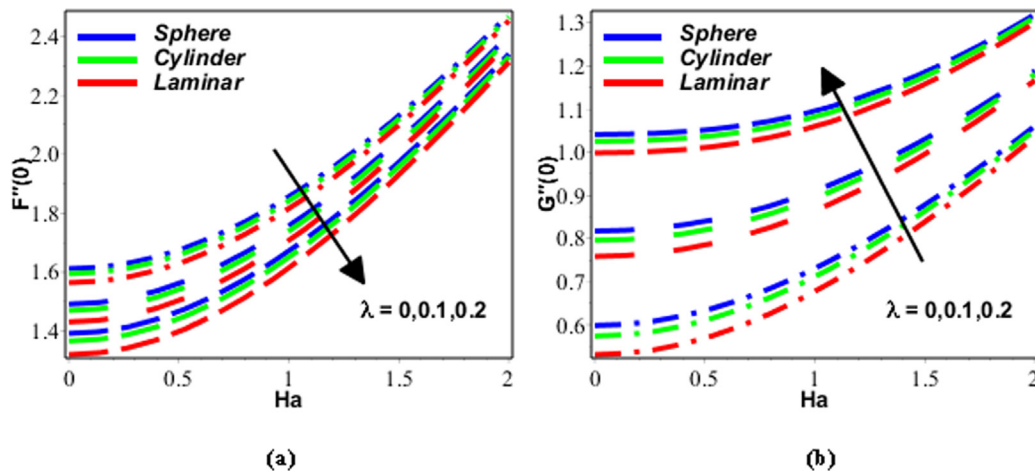


Fig. 2 (a-b). The consequence of  $\lambda$  on (a) skin friction coefficients  $C_{fx}$  (b) skin friction coefficients  $C_{fy}$ .

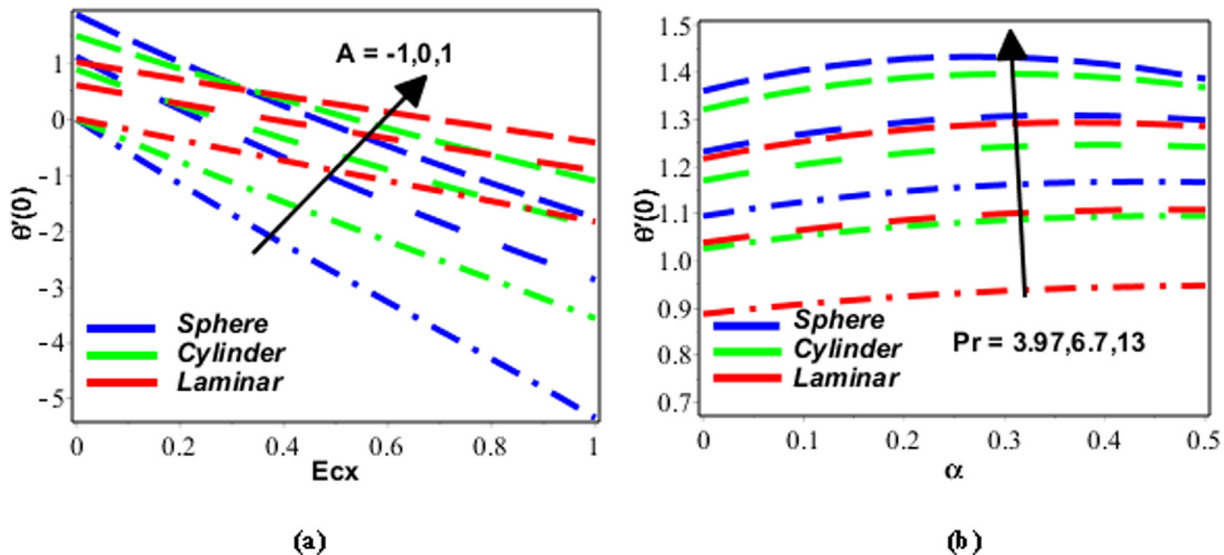


Fig. 3 (a-b). (a) The consequence of  $A$  on Nusselt number ( $Nu_x$ ) (b) Effect of  $Pr$  on Nusselt number ( $Nu_x$ ).

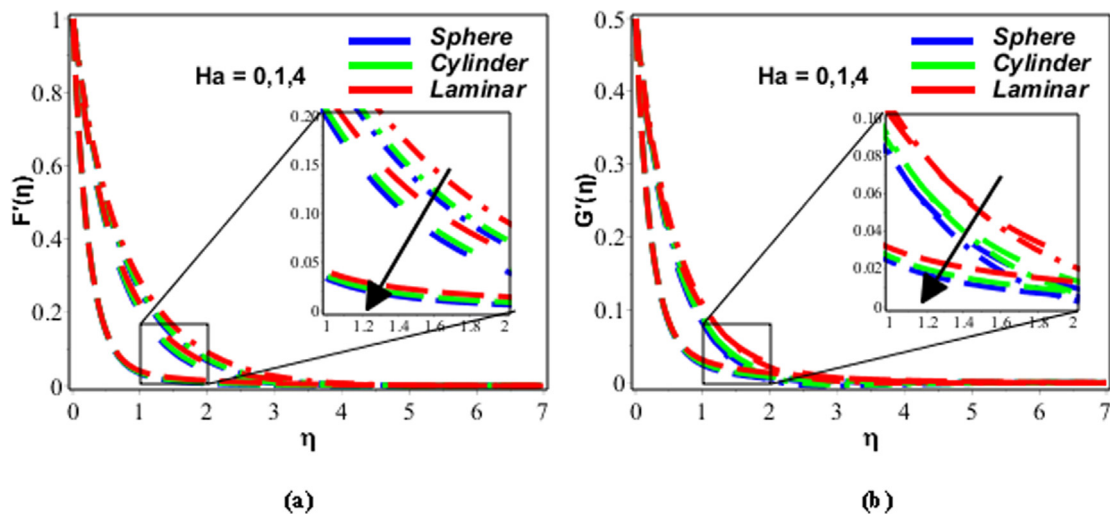


Fig. 4 (a-b). Consequence of  $Ha$  on (a) velocity in  $x$ -direction (b) velocity in  $y$ -direction when  $\lambda = 0.1, \phi = 0.1, Gr = 0.9, \alpha = 0.5$ .

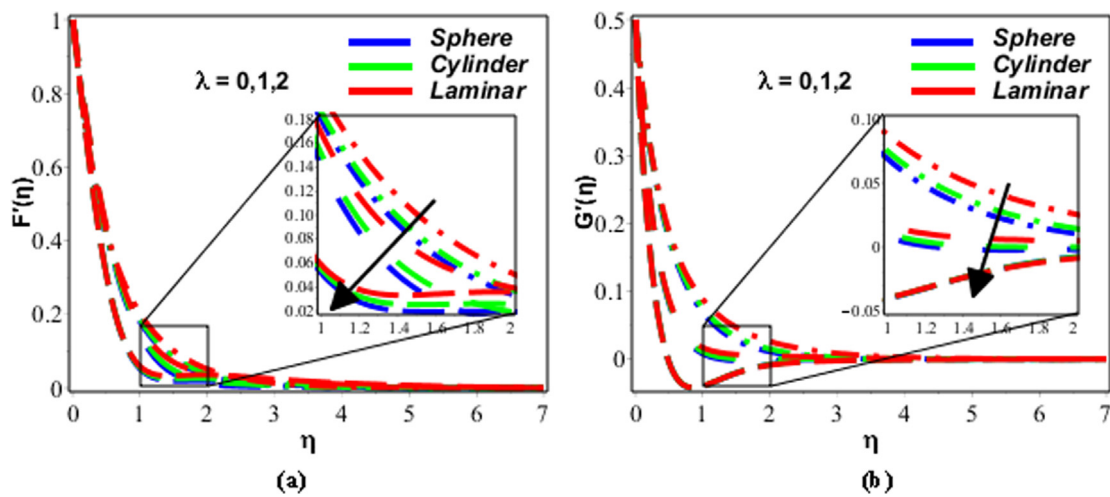


Fig. 5 (a-b). Result of  $\lambda$  on (a) velocity in  $x$ -direction (b) velocity in  $y$ -direction when  $Ha = 2, \phi = 0.1, Gr = 0.9, \alpha = 0.5$ .

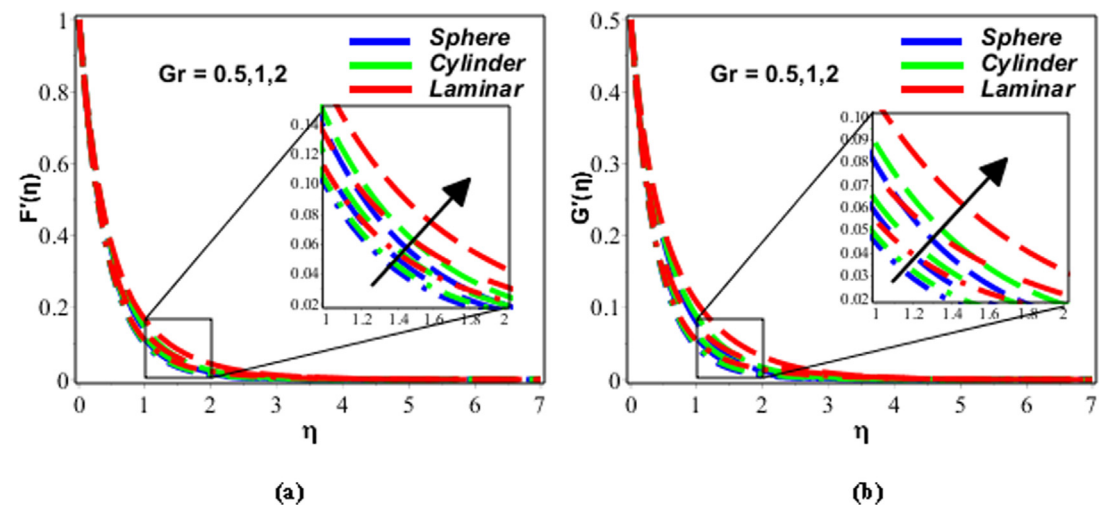
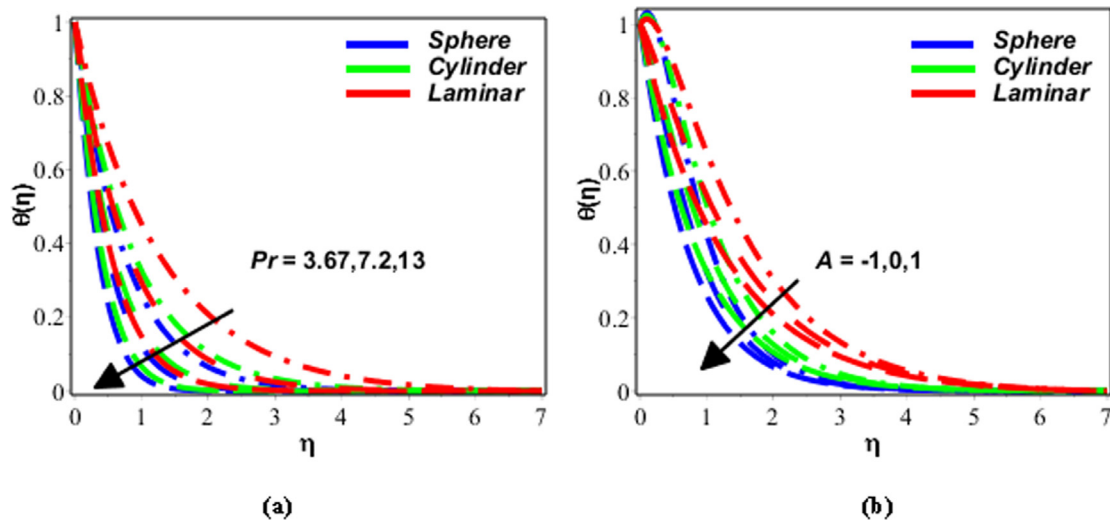
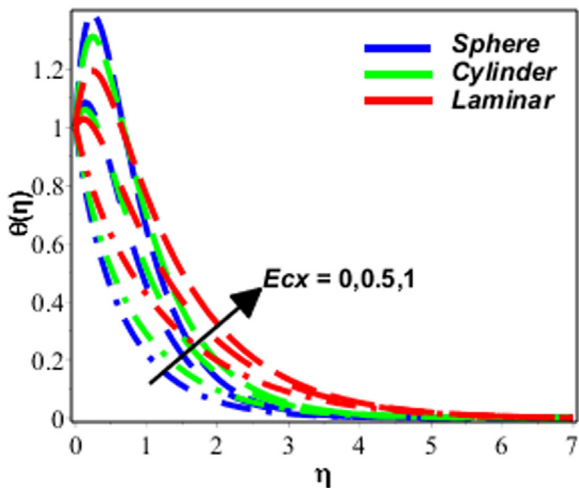


Fig. 6 (a-b). Outcome of  $Gr$  on (a) velocity in  $x$ -direction (b) velocity in  $y$ -direction when  $Ha = 2, \lambda = 0.1, \phi = 0.1, \alpha = 0.5$ .



**Fig. 7** (a-b). Outcome of (a)  $Pr$  on tempertaure for  $Ha = 2, \lambda = 0.9, \phi = 0.1, Gr = 0.9, \alpha = 0.5, Rd = 0.5, Ec_y = 0.5, Ec_x = 0.5, A = 0.9$  (b)  $A$  on tempertaure for  $Ha = 2, \lambda = 0.9, \phi = 0.1, Gr = 0.9, Rd = 0.5, Ec_y = 0.5, Pr = 3.97, Ec_x = 0.3$ .





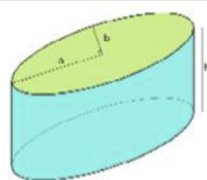
**Fig. 8** Consequence of  $Ec_x$  on tempertaure for  $Ha = 2, \lambda = 0.9, \phi = 0.1, Gr = 0.9, \alpha = 0.5, Rd = 0.5, Ec_y = 0.5, Pr = 3.97, A = 0.9$ .

( $x$ -component) concerning aggregating numerical values of parameter Hartman number is presented. Fig. 4(a) elucidate that the velocity component of the system is lessening with cumulative values of  $Ha^2$ . With the help of the physical study of Hartman number, we can easily claim that this parameter

exists because of the noteworthy magnetic field, which plays the role of an opponent of the fluid to flow. Due to this opposition regarding the flow of fluid, the fluidic velocity is gradually decreasing. Similar effects are observed in Fig. 4(b) for  $y$ -component of velocity. On the other hand, velocity due to the laminar particle is high compared to other particles like spherical and cylindrical shapes. In the subsequent Fig. 5(a-b), effects of velocity ( $x$  and  $y$ -components) due to variated values of rotation parameter are offered, and observations show that velocity is lessening concerning rotation parameter, which is growing. That is why increasing rotation effects have produced the disturbance in the flow of fluid, and hence velocity is increased. On the other side, in comparison between Fig. 4(a-b) and 5(a-b), similar effects regarding the shape of the particles are observed. In Fig. 6(a-d), the performance of velocity with altered  $Gr$  is categorized. Grashof number is reinforced to boost the nano-fluid velocity.

For the sake of observations regarding temperature, Figs. 7 and 8 are strategized. In Fig. 7(a), the declining effects of temperature with increasing values Prandtl number are explained graphically. In the view of physical aspects of Prandtl number, this can easily understand that thermal diffusion of the nano-fluid and Prandtl number is connected with a specific relation named inversely proportional relation. Due to this particular reason, the temperature of the system (based on nano-particles) is increased. In the next, Fig. 7(b) is explained the decreasing effects of temperature due to the enhancement of the temperature exponent parameter. Fig. 8 is strategized to

**Table 3** Different Shapes of nano-particles [29–35].

		
Sphere ( $m = 3$ )	Cylinder ( $m = 6.3698$ )	Laminar ( $m = 16.1576$ )

**Table 4** Comparison of the skin friction values achieved from MCWM with available results in the literature.

Pr	A	$\theta'(0)$		Obtained results
		[38]	[34]	
1	1.5	0.377413	0.37741256	0.37741345
	0	0.549643	0.54964375	0.54964377
	1	0.954782	0.95478270	0.95478269
	3	1.560294	1.56029540	1.56029542
5	1.5	1.353240	1.35324050	1.35324049
	0	1.521243	1.52123900	1.52123898
	1	2.500135	2.50013157	2.50013155
	3	3.886555	3.88655510	3.88655507
10	1.5	2.200000	2.20002816	2.20002819
	0	2.257429	2.25742372	2.25742373
	1	3.660379	3.66037218	3.66037222
	3	5.635369	5.62819631	5.62819628

enlighten the arrogance of temperature for growing numerical values of  $Ec_x$ . The increasing attitude of temperature in the presented Fig. 8, is due to the directly proportional relation between Eckhart number and viscosity of the fluid. Increase in temperature is also due that  $Ec_x$  is the coefficient of viscous dissipation, which corresponds to the enhancement of heat dissipations due to some friction forces. As a result, a certain amount of heat energy is added to the fluid. The change in shape also affects the temperature of the fluid, which can be seen in the Fig. 8 (see Table 3).

Table 4 is presented to assess the important factor “skin friction” accomplished from MCWM with the currently available scheme in refs. [29,33]. Numerical results show the accuracy and compatibility of the suggested scheme. Residual error offered in Table 5 whose expressions are given below as:

**Table 5** Residual error for dissimilar order of estimates.

Order of approximation	$\nabla_F$	$\nabla_G$	$\nabla_\theta$
4	$1.0141 \times 10^{-04}$	$3.0471 \times 10^{-03}$	$3.6187 \times 10^{-05}$
7	$2.1540 \times 10^{-13}$	$4.1004 \times 10^{-10}$	$3.0556 \times 10^{-12}$
12	$3.6540 \times 10^{-20}$	$5.1140 \times 10^{-18}$	$1.5470 \times 10^{-18}$
17	$5.0115 \times 10^{-22}$	$3.1120 \times 10^{-22}$	$5.0580 \times 10^{-23}$
21	$4.7781 \times 10^{-33}$	$3.2512 \times 10^{-31}$	$1.5820 \times 10^{-33}$

modified version of the algorithm to provide better results. Also, the consequences of radiation and  $Gr$  are touched on in the discussion. As a result, the following conclusions are outlined:

$$\nabla_F = \int_0^{\infty} \left[ \frac{\nu_{nf}}{\nu_f} F''' + F''(F+G) - 2F'(F'+G') + 4\lambda G' + \frac{\rho_f}{\rho_{nf}} \frac{(\rho\beta)_{nf}}{(\rho\beta)_f} Gr \theta - \frac{\rho_f}{\rho_{nf}} \frac{\sigma_{nf}}{\sigma_f} Ha^2 F' \right]^2 dq,$$

$$\nabla_G = \int_0^{\infty} \left[ \frac{\nu_{nf}}{\nu_f} G''' + G''(F+G) - 2G'(F'+G') - 4\lambda F' + \frac{\rho_f}{\rho_{nf}} \frac{(\rho\beta)_{nf}}{(\rho\beta)_f} Gr \theta - \frac{\rho_f}{\rho_{nf}} \frac{\sigma_{nf}}{\sigma_f} Ha^2 G' \right]^2 dq,$$

$$\nabla_\theta = \int_0^{\infty} \left[ \frac{1}{Pr} \left( \frac{\kappa_{nf}}{\kappa_f} + Rd \right) \theta'' + \frac{(\rho c_p)_{nf}}{(\rho c_p)_f} \left( (F+G)\theta' - A(F'+G')\theta \right) + Ec_x Ec_y (F'^2 + G'^2) \right]^2 dq.$$

It can be observed that error  $[\nabla_F, \nabla_G, \nabla_\theta] \rightarrow 0$  as enhancing the order of approximations.

**6. Conclusion**

This report covers both the various ways heat is transferred, as well as forms of  $Cu$ -nanoparticles, shown here. We use a

- The modified Chebyshev wavelets technique lowers computing effort and is very efficient in solving problems.
- Nanofluids based on copper oxide nano-particles have slow velocity than the other nano-fluids. This is due to the copper oxide nano-particles, which enhance the viscosity of the nano-fluid over the different nano-fluids. It allows the nano-fluid to take more time on the heated surface and thus absorbs more heat from the surface than other nano-fluids.



- “Skin friction” demonstrates the increasing behavior against the Hartmann number. Due to the variation in rotating parameter,  $x$ -directional “skin friction” decreases gradually, and the reverse effect of rotating parameter is achieved for  $y$ -directional skin friction coefficient.
- “Skin friction” has the maximum values for the sphere shape of the nano-particle.
- The Nusselt number increases owing to  $Pr$ , the stretching ratio parameter, and the temperature exponent parameter and decreases as the Eckert number increases.
- The laminar shape of nano-particles is dominated by velocity and temperature profiles.

### Declaration of Competing Interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

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